



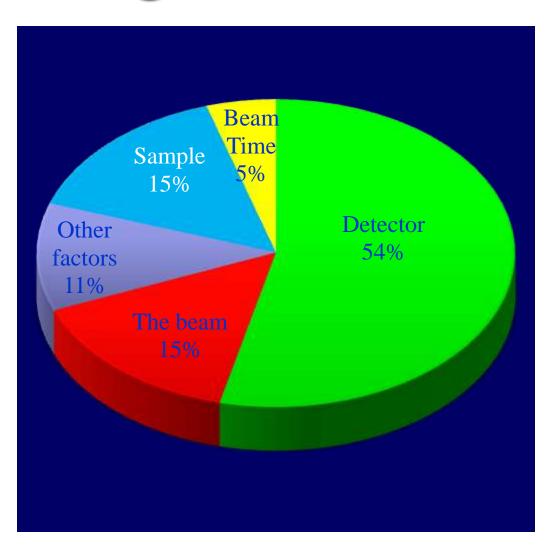
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Detectors for Synchrotron Radiation

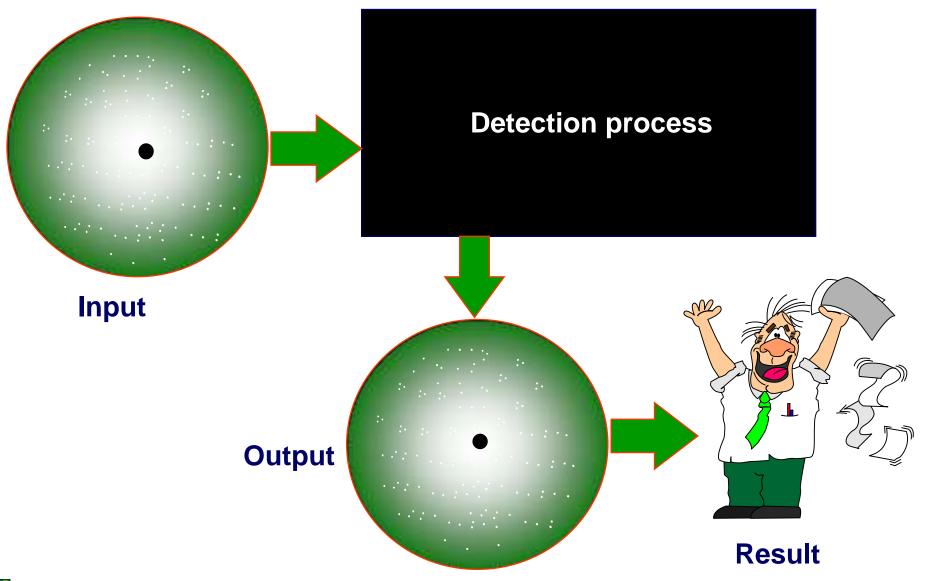
Rob Lewis

Factors Limiting Science

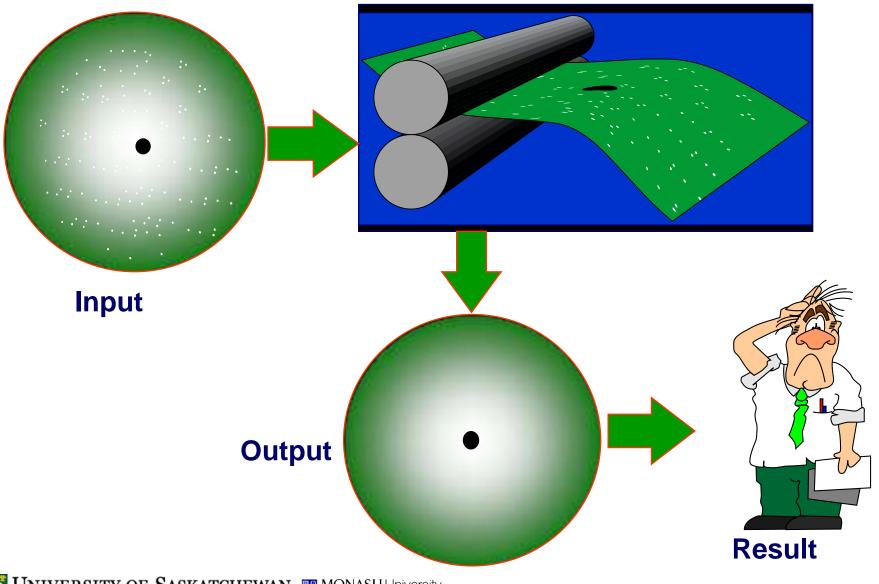
- Detectors are an oftneglected but crucial part of an experiment
 They often limit the
 - science



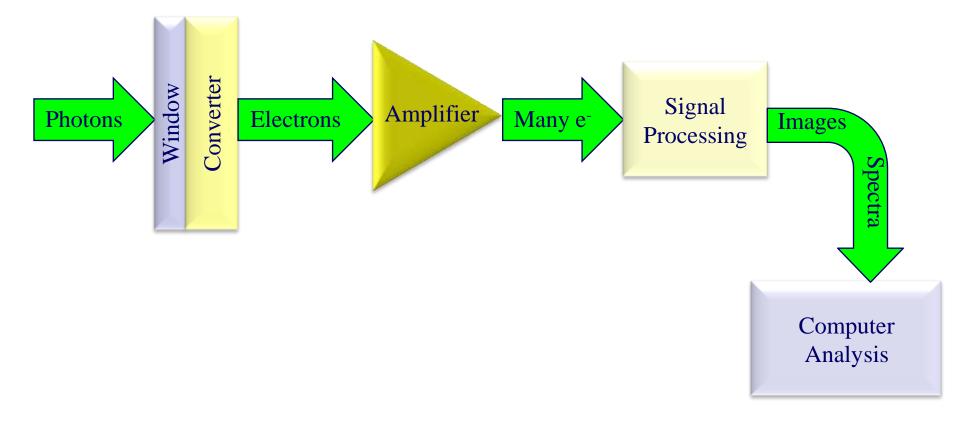
Scientist's View of Detector



The Truth!



Detector Chain of Events



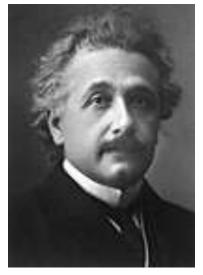
Detection Mechanisms

There are many means of detection. All require the interaction of photons/electrons with matter

Examples include

- ♦ Gas ionisation
 - Photons produce electrons and ions which are then detected
 - E.g. Ion chambers, proportional counters
- Photoelectric effect
 - Photons eject electrons from a solid creating a current which is measured
 - E.g.. Beam monitors
- Generation of electron hole pairs
 - Photons produce electrons and holes in a semiconductor which are then detected
 - E.g.. CCD
- Fluorescence, scintillation and F centres
 - Photons produce prompt fluorescence or F centres
 - E.g. Image plates and Scintillation counters
- Chemical effect
 - Photons create a chemical change such as dissociating Ag halide
 - E.g. Film

Albert Einstein



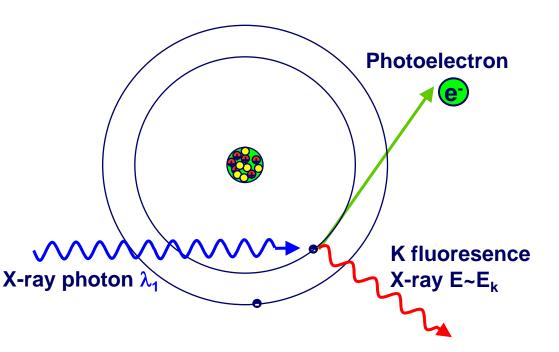
Germany and Switzerland Kaiser-Wilhelm-Institut (now Max-Planck-Institut) für Physik Berlin-Dahlem, Germany **1879 - 1955**



Nobel prize in physics 1921

"for his services to Theoretical Physics, and especially for his discovery of the law of the photoelectric effect"

Photoelectric Effect



Arthur Holly Compton



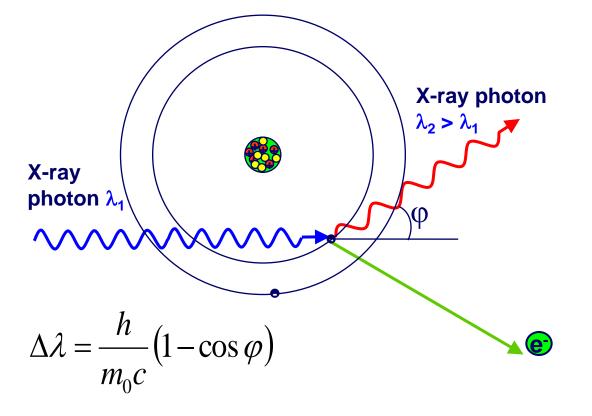
Nobel prize in physics 1927

"for his discovery of the effect named after him"



University of Chicago Chicago, IL, USA **1892 - 1962**

Compton Effect

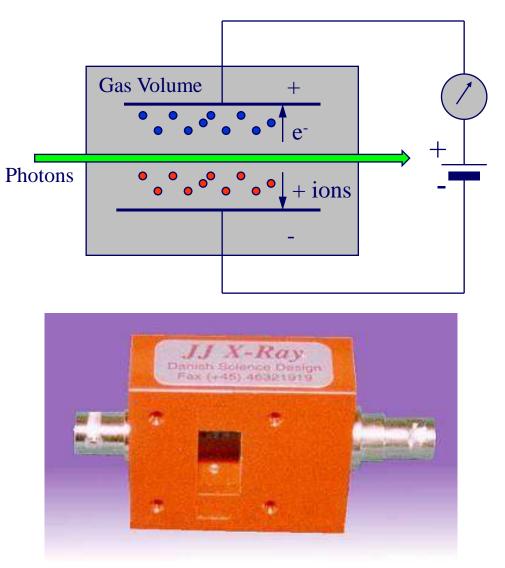


An Example Detector



Echidna

Ionisation Chamber

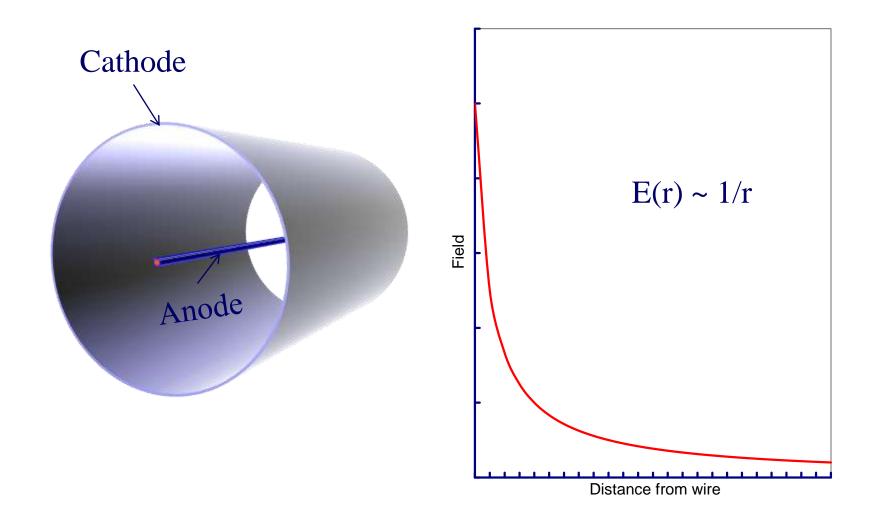


Very simple device

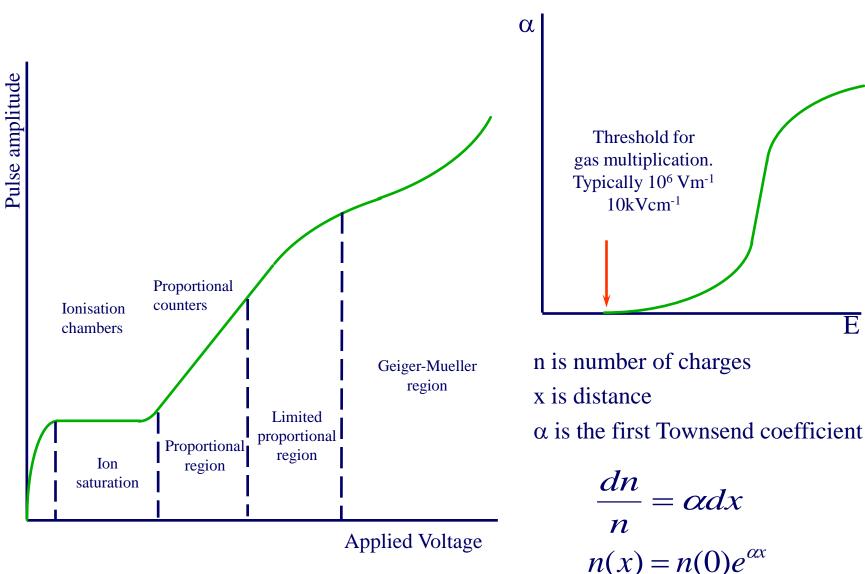
- Approximately 1 e⁻ ion pair per 30eV deposited
- Important that recombination low as possible
 - Higher voltages required at higher rates since more carriers
 - Diffusion losses caused by separation of carriers minimised by higher voltages
 - Plates too close cause electron losses

Ion chambers are sensitive to pressure and temperature

Field Variation

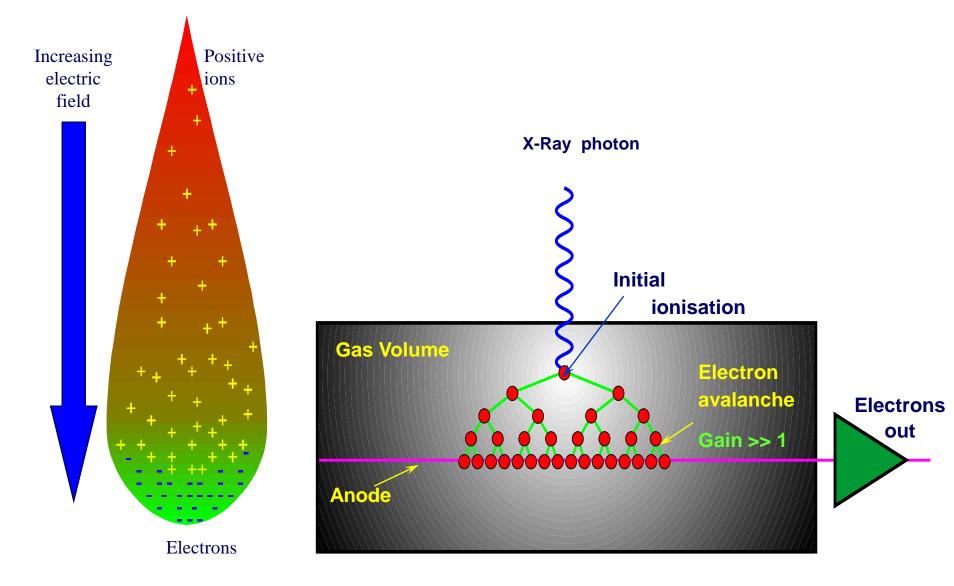


Operation regions of gas filled detectors



E

Avalanche & Proportional Counter



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Georges Charpak



France École Supérieure de Physique et Chimie Paris, France; CERN Geneva, Switzerland

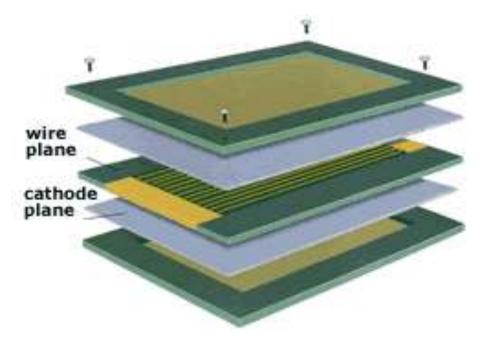
b. 1924 (in Dabrovica, Poland)



Nobel prize in physics 1992

"for his invention and development of particle detectors, in particular the multiwire proportional chamber"

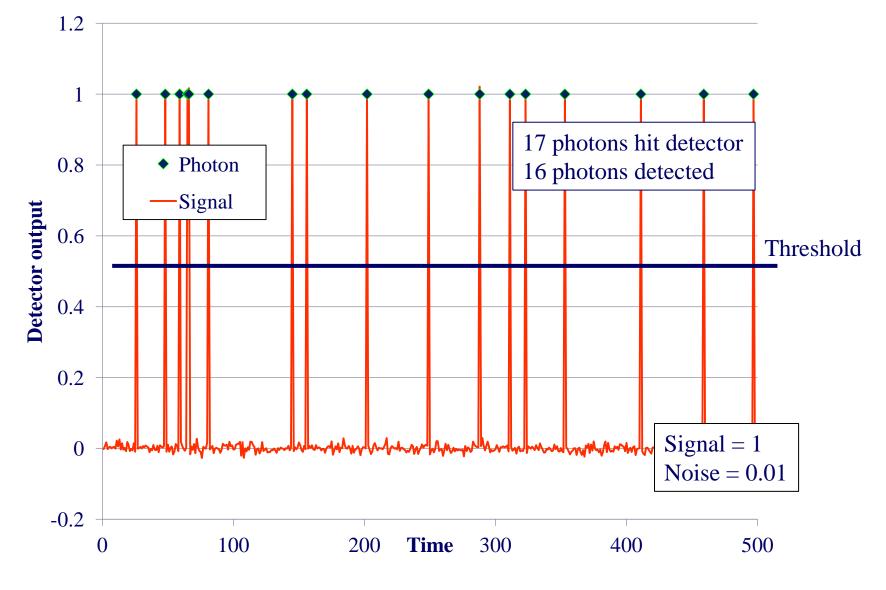
Multi-wire Proportional Counter



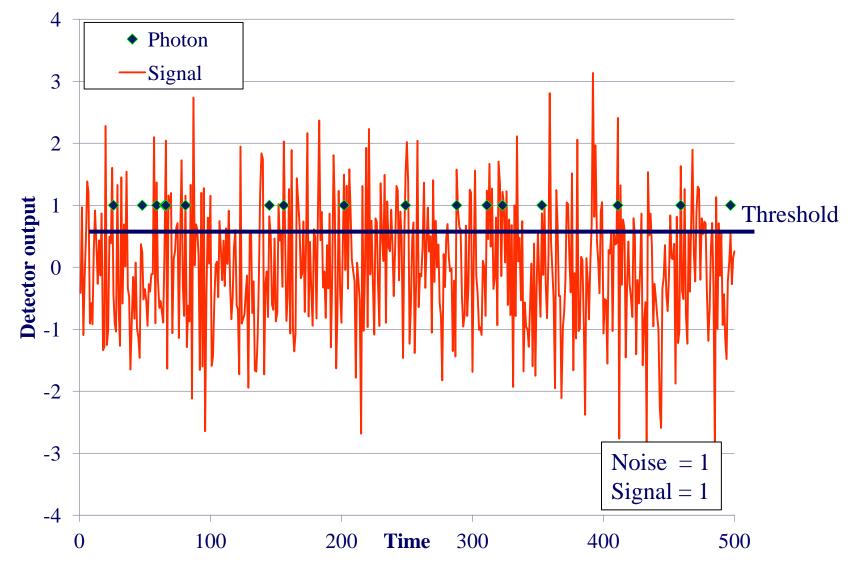
Counting and Integrating

- If there is sufficient signal produced by the interaction of a photon or a particle in the detector then it is possible to operate the detector as a counter
- It's all about signal to noise ratio!

SNR = 100



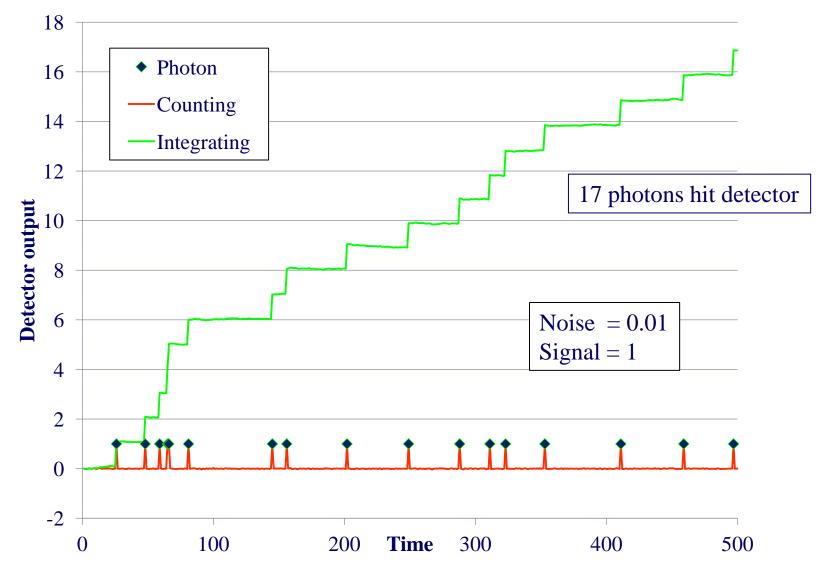
SNR = 1



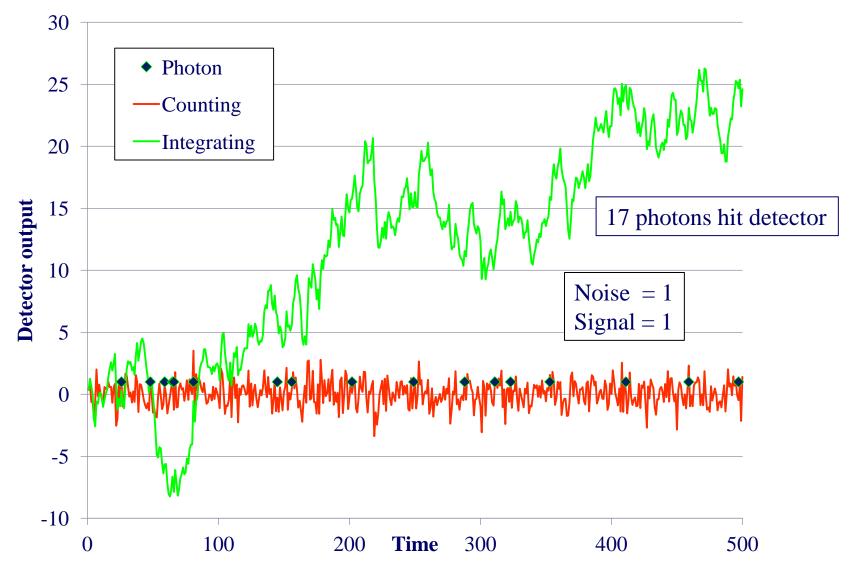
Counting and Integrating

Usually SNR is insufficient and we have to accumulate many photons/particles before the signal becomes measurable

Counting & Integrating SNR =100



Counting & Integrating SNR = 1



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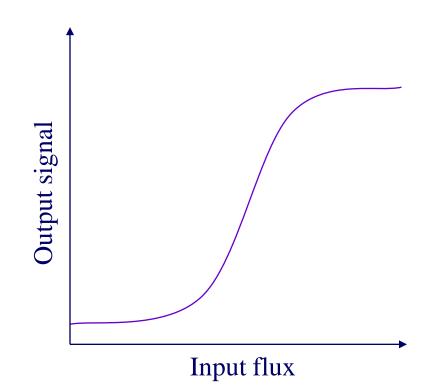
Integrating Detectors

Mode

 Measures deposited energy at end of integration period

Characteristics

- High input flux capability
- Read noise dominates at low signal ("fog level")
- Dead time between frames
- 2×20 keV phts = 1×40 keV photon i.e. Cannot perform simultaneous spectroscopy and positioning
- Examples: Image plates, CCDs



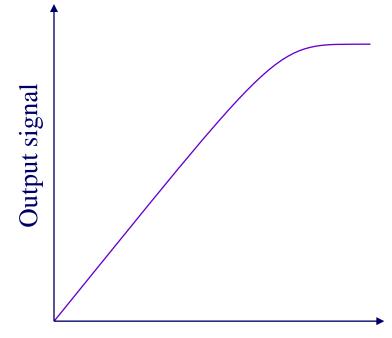
Photon Counting Detectors

Mode

 Detects every photon as it arrives. Only active pixels read

Characteristics

- Quantum limited, Detector noise often negligible
- No dead time between frames
- Can measure position and energy simultaneously
- Limited input flux capability
- Examples: Prop counters, Scintillators



Input flux

Types of Detectors



Crimson Rosella and King Parrot

Willard S. Boyle & George E. Smith

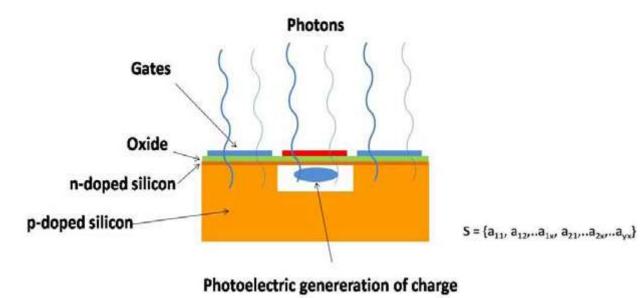


Nobel prize in physics 2009

"for the invention of an imaging semiconductor circuit – the CCD sensor"

123

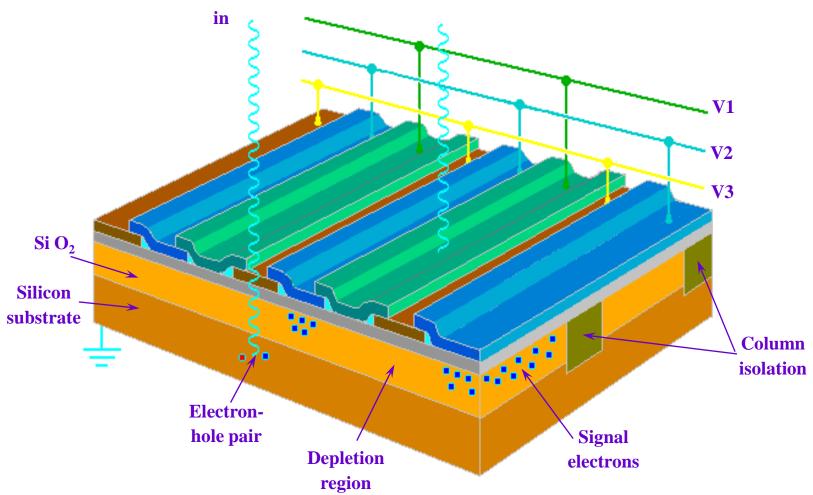
Bell Laboratories Murray Hill, NJ, USA



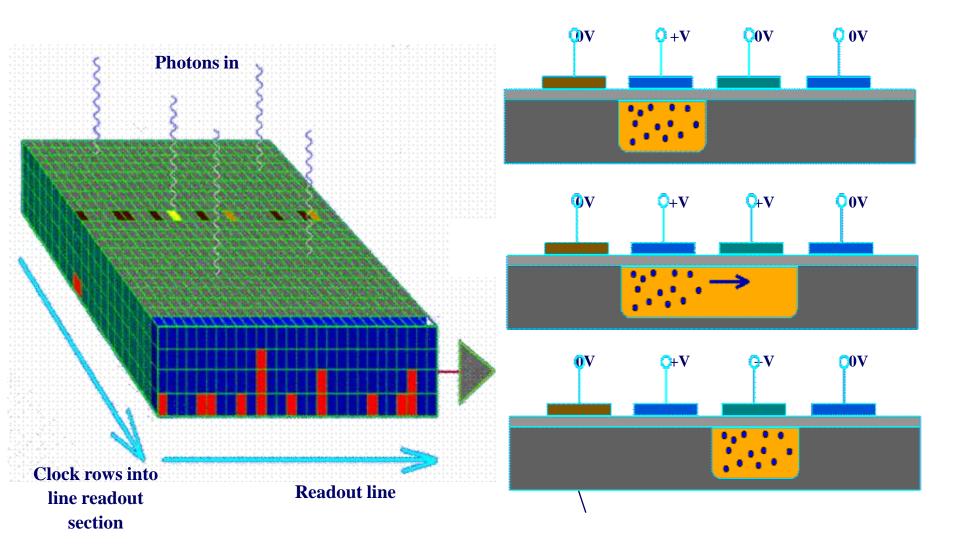
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Charge Coupled Device





CCD Readout

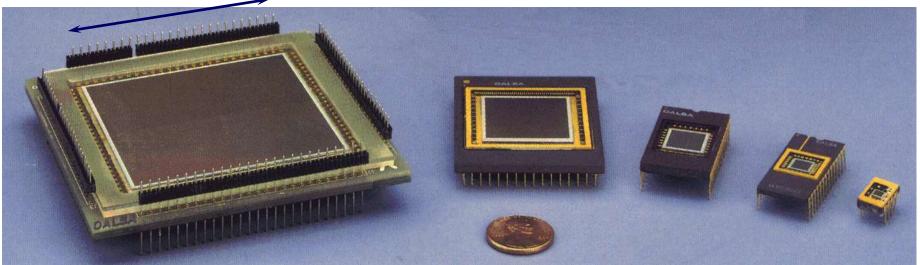


CCD Readout

- Charge is moved from pixel to pixel by clocking
- Each pixel has a limited capacitance (well depth) typically 10⁴-10⁵ e⁻
- This limits dynamic range for direct detection
 - 10keV photon creates ~ $3000e^{-}$ so saturation = ~ 10 photons
- Speed of clocking is restricted by line capacitance and charge transfer efficiency
 - Size of CCD restricted by this
- Noise can be reduced by cooling
- Amplifier usually on chip
 - Heats up that part of chip

CCDs

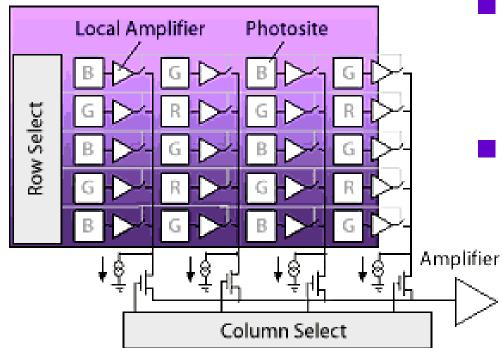
62mm



Although sizes > 50mm are available, the read speed is slow to preserve low noise and cte (line capacitance becomes very high) Shutter required

Complimentary Metal-Oxide Semiconductor (CMOS)

CMOS Imager

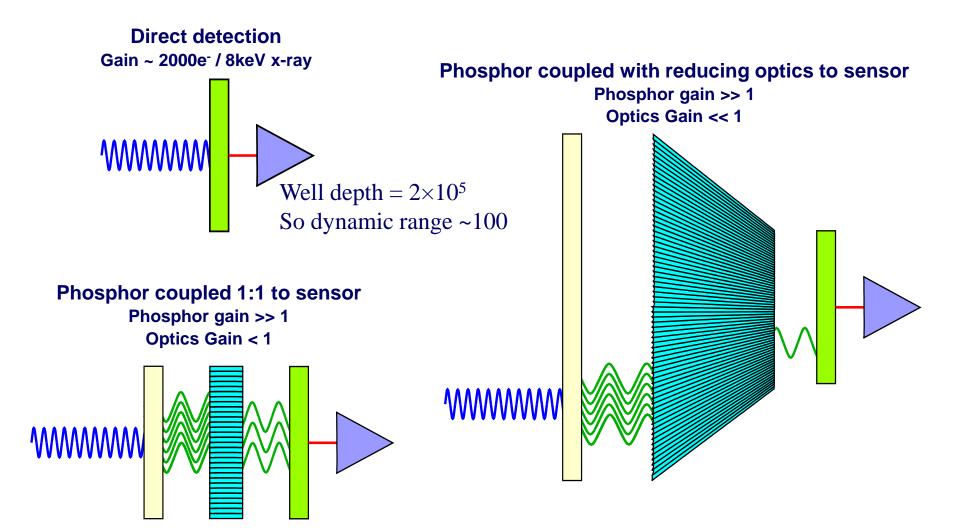


 A readout amplifier transistor on each pixel converts charge to voltage
 Allows random access to pixels, similar to the rowcolumn memory cell access in RAM

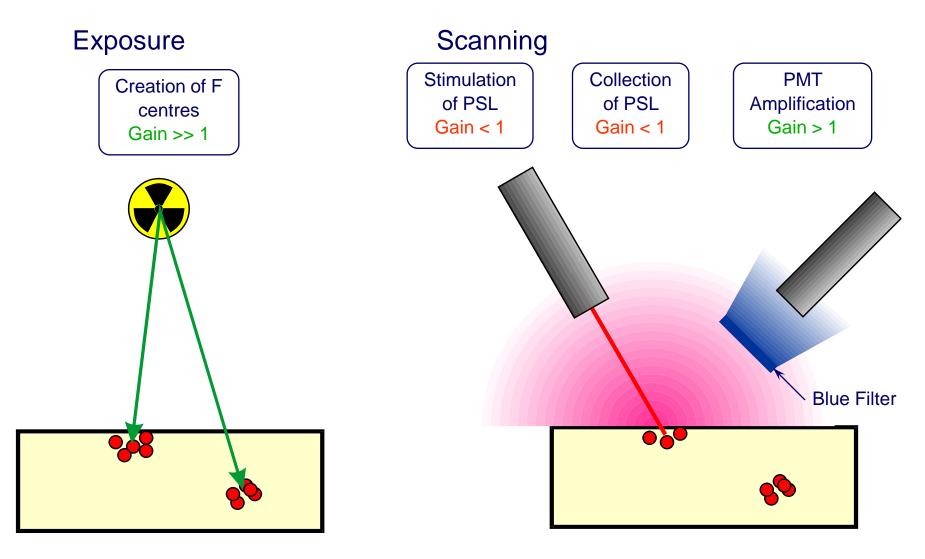
CMOS vs CCD

- Traditionally CCD higher sensitivity and lower noise
- Modern lithography means they are now similar
- CMOS sensors can have much more functionality on-chip than CCDs
 - On chip image processing, edge detection, noise reduction, and analog to digital conversion
- CMOS lower power \rightarrow less heat \rightarrow less noise

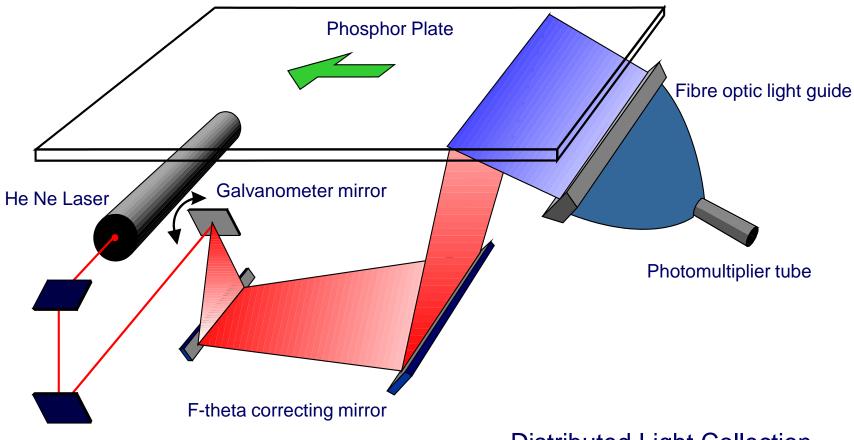
Use with X-rays



Computed Radiography-Image Plate



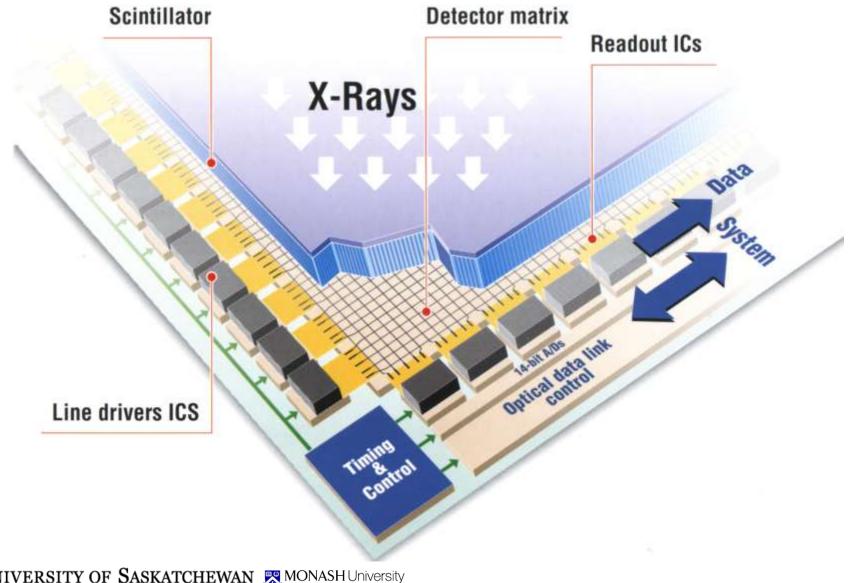
X-Y Flat bed Scanner



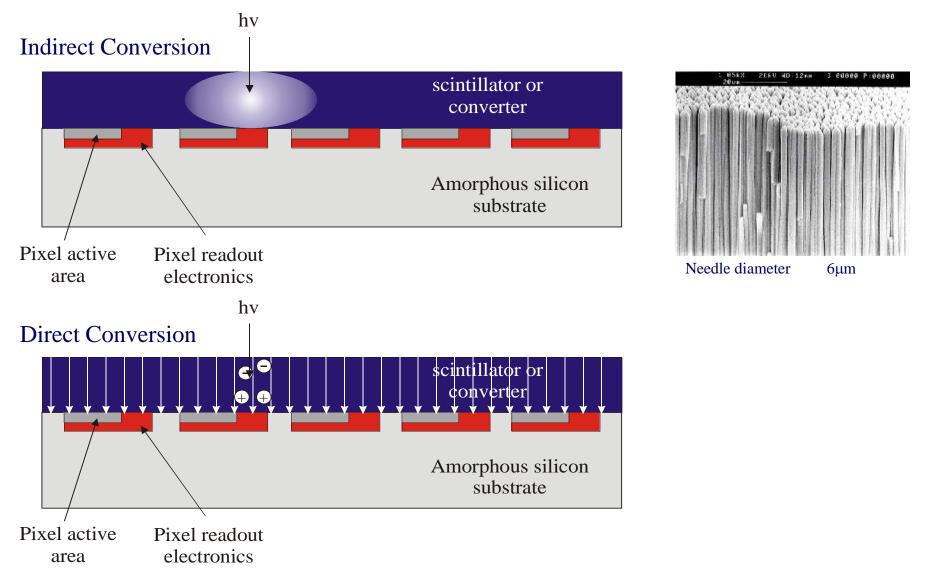
Distributed Light Collection



TFT Flat panel Detector



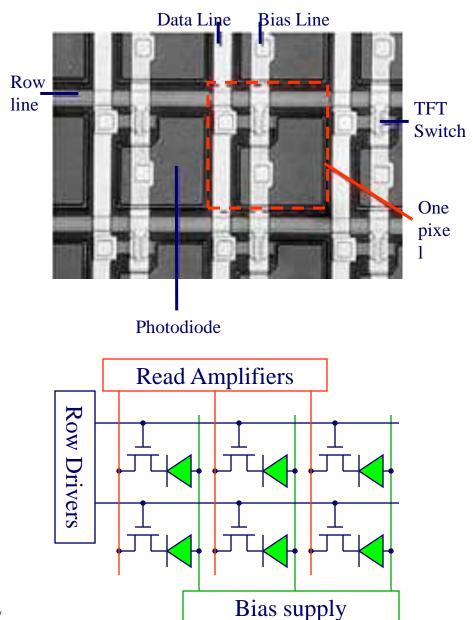
a-Si:H TFT arrays



a-Si:H Array dpiX - Flashscan 30

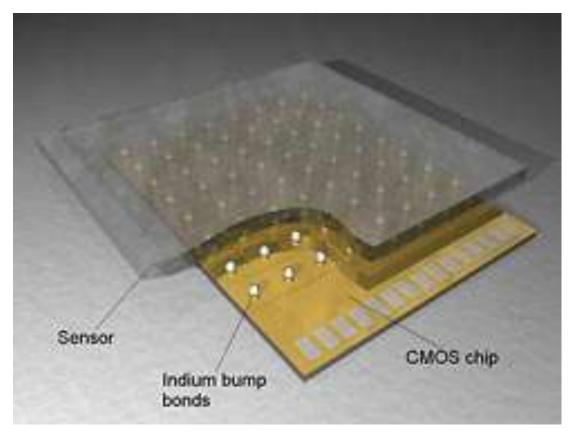


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PILATUS 6M Detector









Ch. Brönnimann, E. Eikenberry, B.Schmitt, M. Naef, G. Hülsen (SLS); R. Horisberger, S. Streuli (TEM); Ch. Buehler (LOG); F. Glaus (LMN); M. Horisberger (LNS)

PILATUS 6M Detector



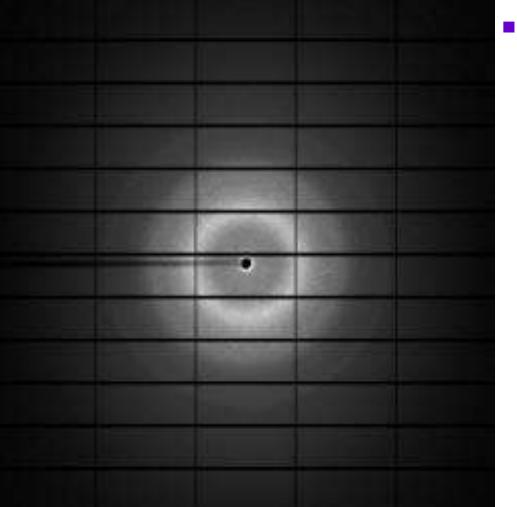


- Reverse-biased silicon diode array
- Thickness 320 μm
- Pixel size 172 x 172 μ m²
- **2463 x 2527 = 6,224,001 pixels**
- Area 431 x 448 mm²
- Intermodule gap x: 7 pixels, y: 17 pixels,
 8.4% of total area
- Dynamic range 20 bits (1:1,048,576)
- Counting rate per pixel > $2 \times 10^6 \text{ X-ray/s}$
- Energy range 3 30 keV
- Quantum efficiency (calculated)
 - 3 keV: 80% 8 keV: 99% 15 keV: 55%
- Energy resolution 500 eV
- Adjustable threshold range 2 20 keV Threshold dispersion 50 eV
- Readout time 3.6 ms
- Framing rate 12 Hz
- Point-spread function 1 pixel

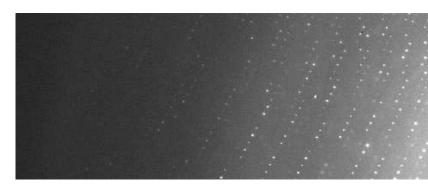
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PILATUS 6M Detector





 X-ray diffraction image recorded from a ferritin crystal (energy=16 keV, distance = 204 mm).



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Spectroscopic Detectors



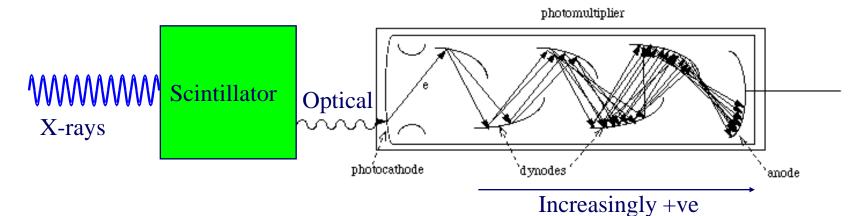
Rainbow Lorikeets



Spectroscopic Detectors

- For quantitative work, most are counting detectors that measure the size of individual energy deposits
- Alternative ids the use of filters as in optical colour cameras

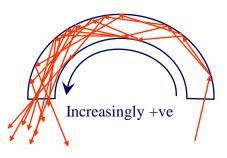
Electron multipliers & Scintillators





Channeltron is a similar with distributed dynode

Micro-channel plates are mutlichannel channeltrons with each channel being an electron multiplier.



Multi Channel Spectoscopic Detectors





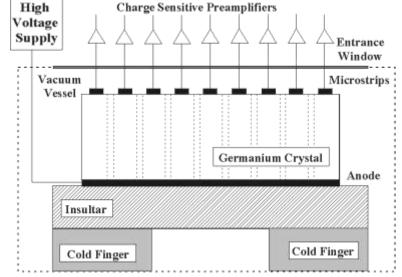
Canberra Ultra-LEGe detector

WRULEAD (Windowless, Retractable, Ultra Low Energy Array Detector) works down to 300eV

Multichannel devices up to 30 channels at 3×10^5 cts s⁻¹ channel⁻¹ have been built

SPring-8 128 channel Ge strip





Ge ◆ 55.5×50.5×6mm **Strips** Number 128 300µm • Width • Interstrip 50µm Length 5mm Readout Single channel 100ns ♦ 32 channels 3.2ms Max expected count rate ♦ 14kcps

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Spectral Resolution

Average number of carriers, N = E/w where w is energy to create electron hole/ion pair
 Poisson statistics σ = 1/√N

$$= (E/w)^{-1/2} = (w/E)^{1/2}$$

• $\Delta E/E \text{ fwhm} = 2.355\sigma$ = $2.355(w/E)^{\frac{1}{2}}$

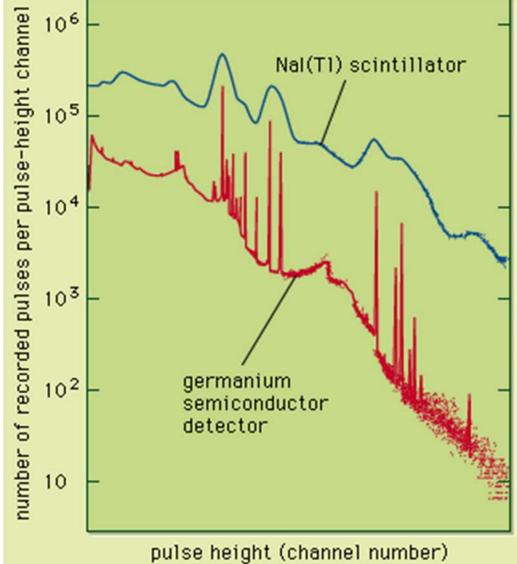
For Ge, w = 3eV so at 10keV ΔE/E ~ 4%
For NaI, w = 30eV so at 10keV ΔE/E ~ 13%

Fano Factor

- If all energy from photon or particle were converted into carriers there would be no variance
- Poisson statistics assume only a small fraction of energy goes into charge creation
- Reality is somewhere in between so we introduce Fano factor F
- Fano factor is defined as $F = \frac{\sigma^2}{\mu}$ where σ^2 is the variance and μ is the mean number of carriers
- For a Poisson process, the variance equals the mean, so F = 1
- Examples
 - Si: 0.115
 Ge: 0.13
 GaAs: 0.10
 Diamond: 0.08

• Observed relative variance = F x Poisson relative variance

Scintillator vs Germanium



The top spectrum is from a scintillation detector, and the bottom is from a germanium semiconductor detector. The superior energy resolution of the germanium is evident from the much narrower peaks, allowing separation of gamma-ray energies that are unresolved in the scintillator spectrum.

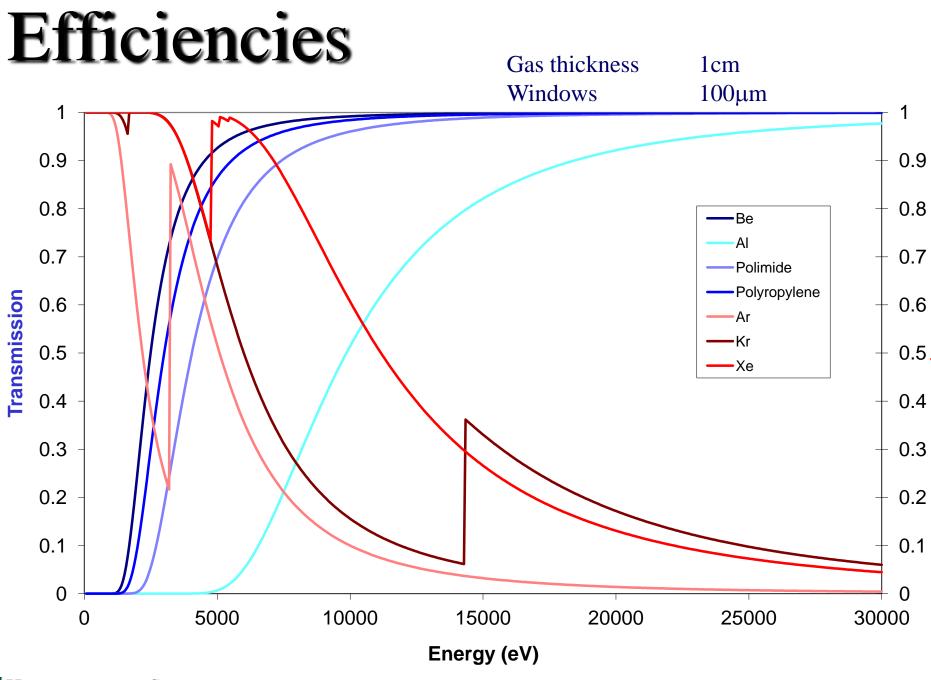
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Things to Look Out For



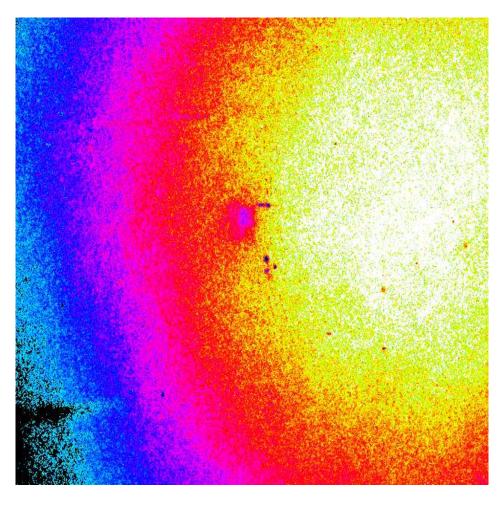
Crocodile



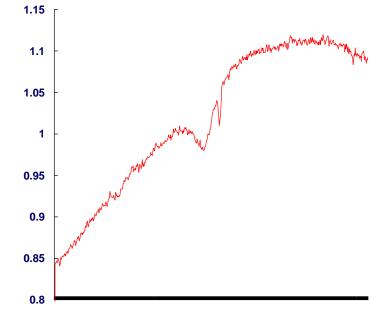


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Response to Uniform Illumination

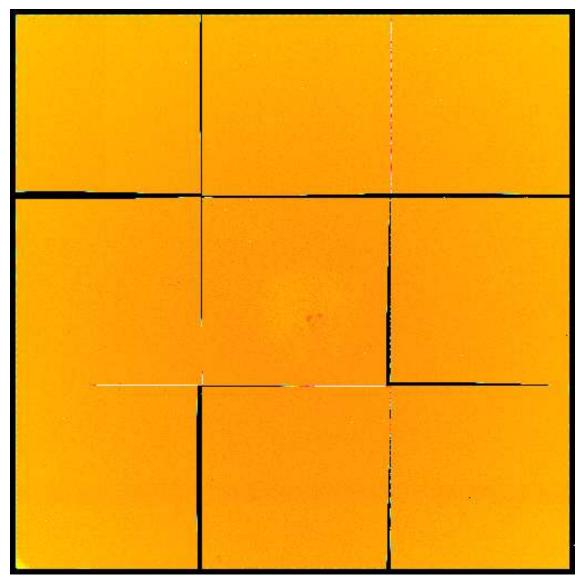


ESRF TV Detector Thompson IIT & CCD



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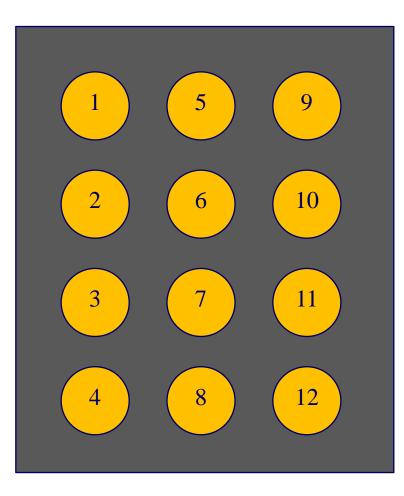
Gaps



Spec	0.2mm max
Worst gap	2.97mm
Pixels in gaps	513922 5.45%

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Intensity Test



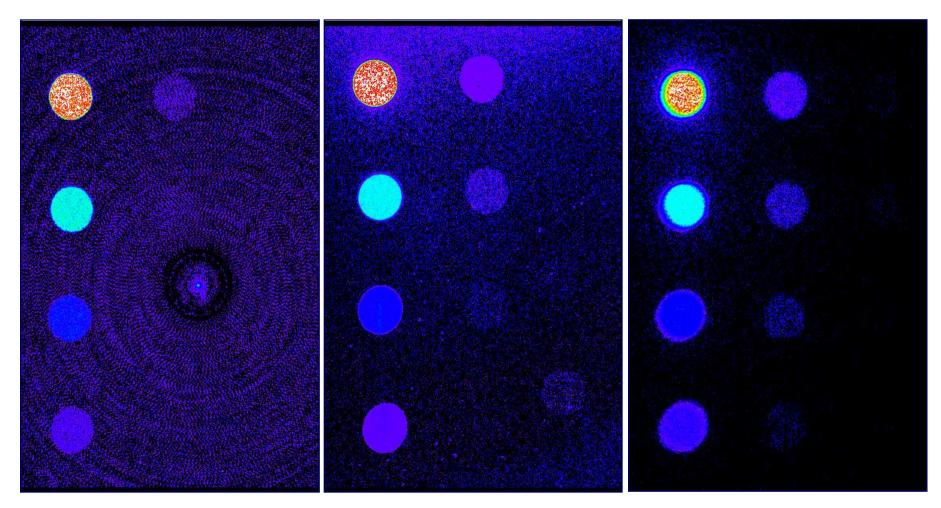
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Graded Absorber Comparison

Mar Image Plate

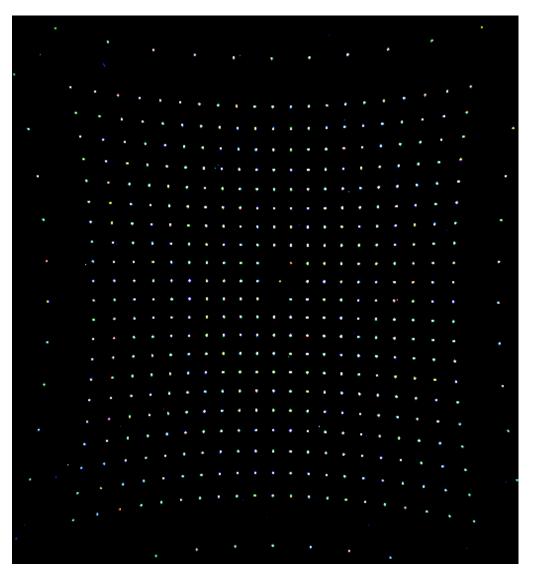
ESRF-Thompson IIT / CCD

Daresbury MWPC



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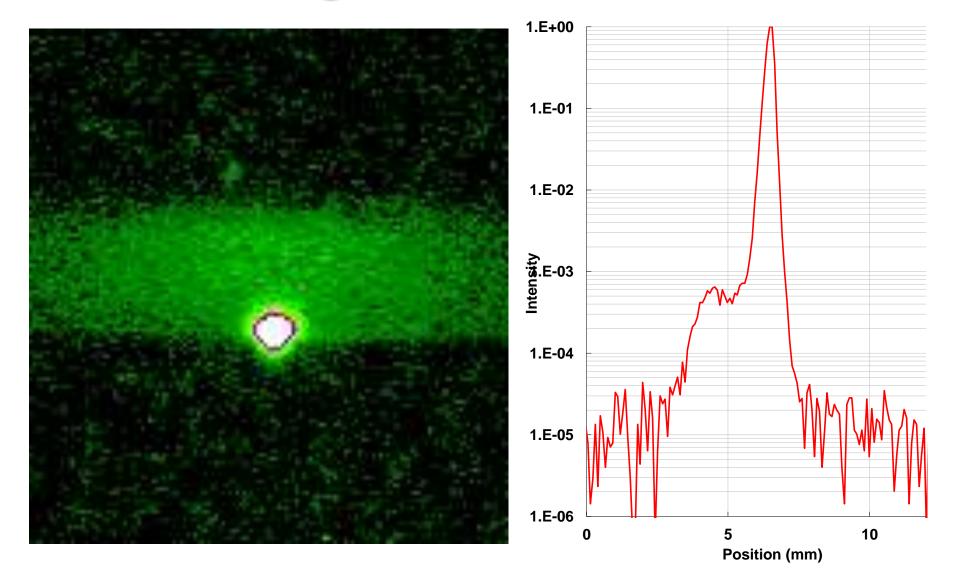
Spatial distortion



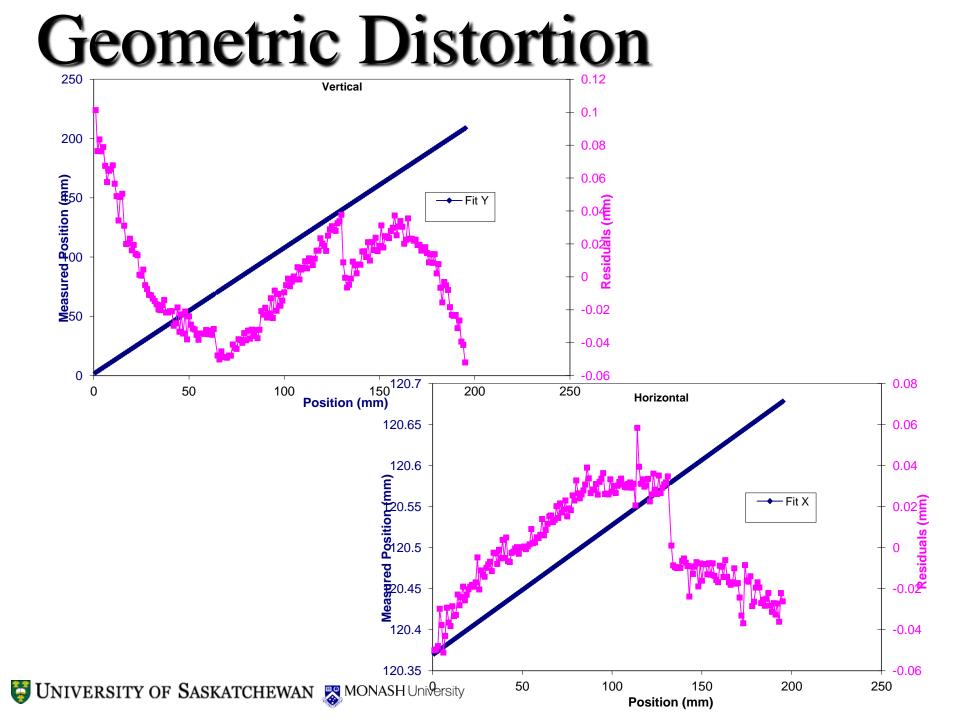
ESRF Image intensifier detector

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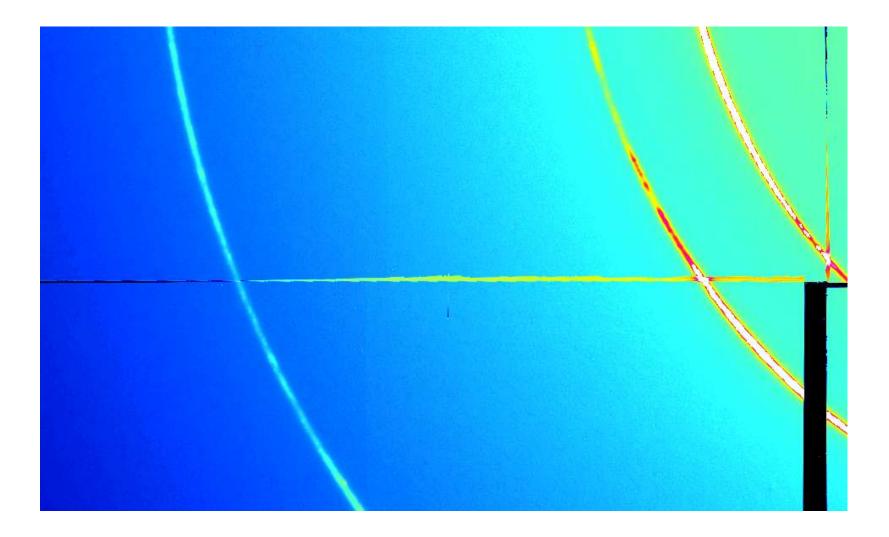
IPlate Single Peak PSF



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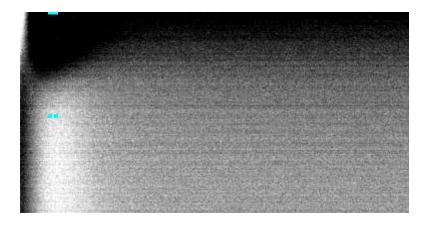


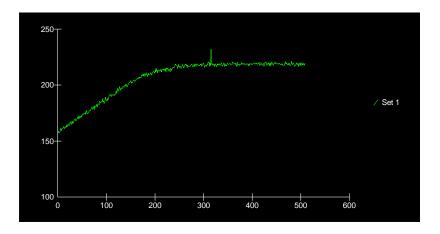
Overlaps



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Dark Currents





Flat and Dark Correction

For each image, two correction images must be recorded.

1. A flat field (uniform illumination of the detector)

2. A dark image (no irradiation of detector)

Both must be recorded with the same exposure time as the original image since dark current is a function of exposure time.

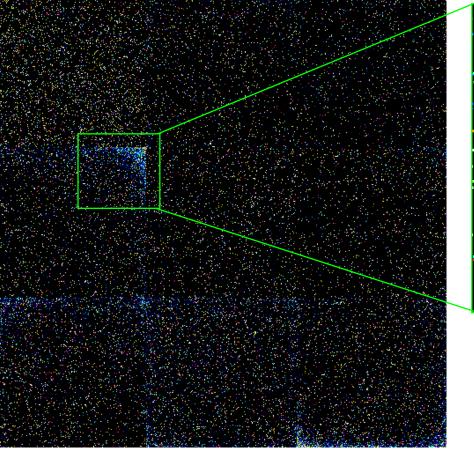
Then apply the following correction

$$Corrected = \frac{(image - dark)}{(flat - dark)}$$

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Dark Current

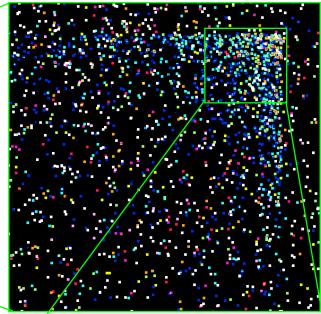
Pixels above the 0.2 photons pix⁻¹ specification

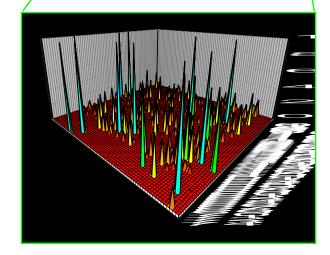


Number failing 2 measurements 5-2000s

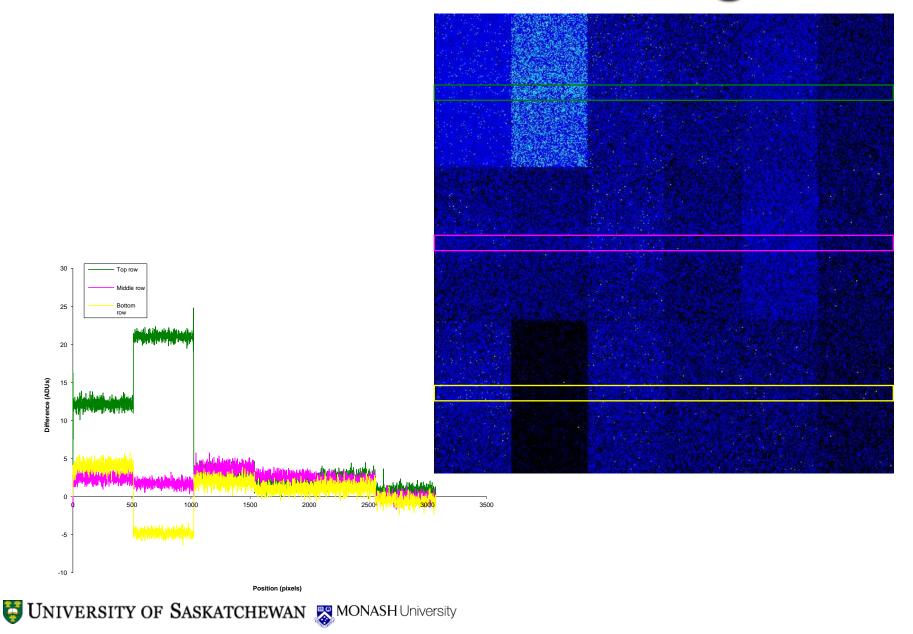
Mean	44764	0.47%
Min	40822	0.43%
Max	48706	0.52%
nb. 14300 pixels not common to both		

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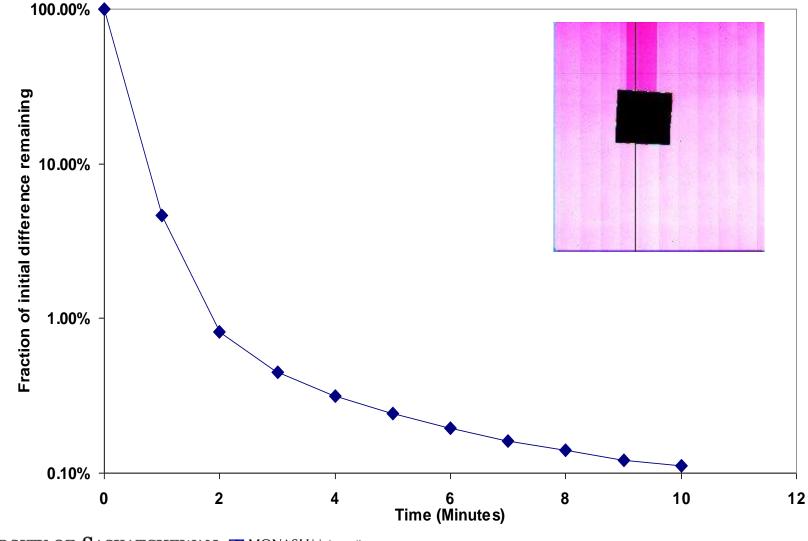




Subtraction of dark images



Flashscan 30 - Image Lag



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Radiation Damage (Medipix)

- Damage occurred at 40Gy or 1.3×10¹⁰pht/mm² in the readout chip
- At 13 keV photon energy
 - Strong diffraction spots typically 10⁵ phts/s or 10⁶ phts/mm2/s
 - Damage requires ~ 8hours exposure
 - Direct beam (10¹⁰–10¹³ photons/mm²/s)
 - Damage in less than a second.

dpiX Flashscan 30 PaxScan 4030

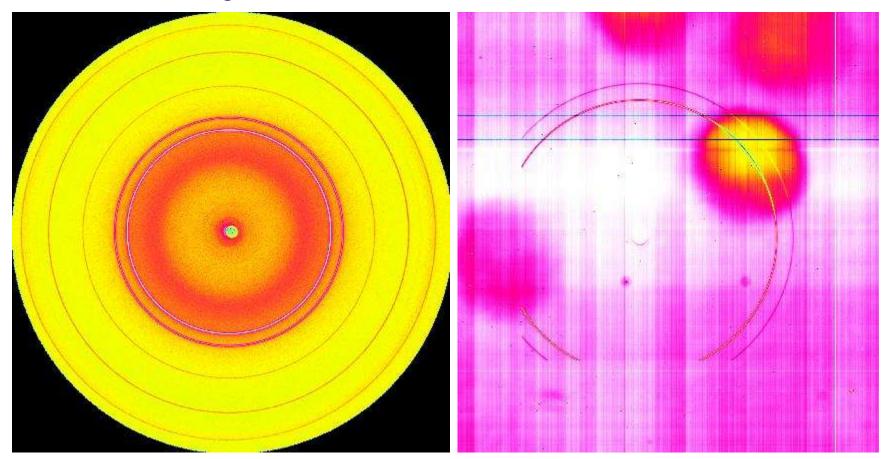


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Flashscan 30 - Performance

Mar Image Plate

Flashscan-30



 $t_{int} = 30s$

t_{int}=190s

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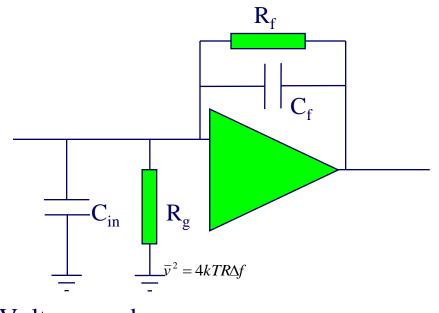
Electronics Issues



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Albino Kookaburra

Amplification



- Voltage mode
 - Output \propto input voltage
 - Effect of R_f dominates C_f
- Current mode
 - ♦ Output ∝ input current
 - Low input impedance
- Charge mode
 - Output ∞ input charge
 - C_f dominates R_f

- In almost all cases we require amplification
- Amplifier-detector interaction is critical
- Most important element is the input
- Noise is the major issue
 - Thermal or Johnson Noise
 - Brownian motion of electrons
 - No current flow or voltage required
 - White noise
 - Shot Noise
 - Fluctuations in current
 - White noise

$$\bar{i}^2 = 2q_e\bar{I}\Delta f$$

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Equivalent Noise Charge

Introduce ENC which is that signal charge that will produce the same output as the RMS noise

$$ENC^{2} = \exp\left(2\right)\left[\frac{kT}{2R_{g}}\tau + \frac{eI_{D}}{4}\tau + \frac{kT(C_{in})^{2}}{2g_{m}\tau}\right]$$

Where

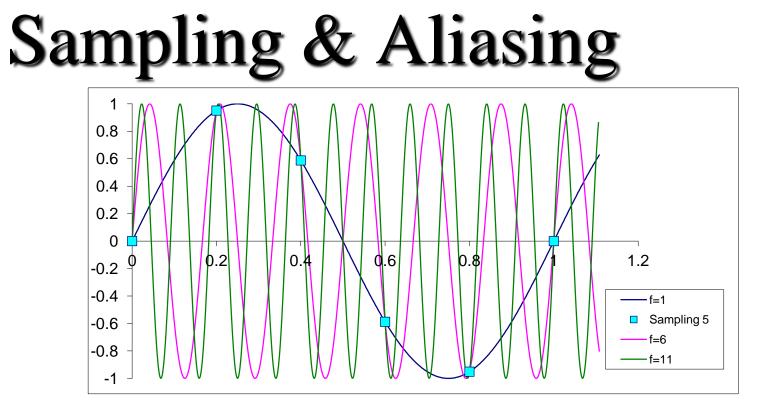
- k = Boltzman's constant
- T = temperature
- e = the electronic charge
- \mathbf{R}_{g} = Load resistance and/or feedback resistance
- $g_m = \text{transconductance of input FET. (Links current in to voltage out)}$
- $\tau = \text{Rise time of amplifier}$
 - C_{in} = input / stray and feedback capacitance
 - I_D = Drain current
- Note that ENC is directly related to energy resolution
- FWHM(keV) = 2.355×10^{-3} ENC/ew where w is the energy per electron

Noise Dependence $ENC^{2} = e^{2} \left[\frac{kT}{2R_{f}} \tau + \frac{q_{e}I_{D}}{4} \tau + \frac{kT(C_{in})^{2}}{2g_{m}\tau} \right]$

• τ optimum at

$$\tau_{opt} = \left[\frac{kT/2g_m}{(kT/2R_f) + (q_e I_D/4)}\right]^2 C_{in}$$

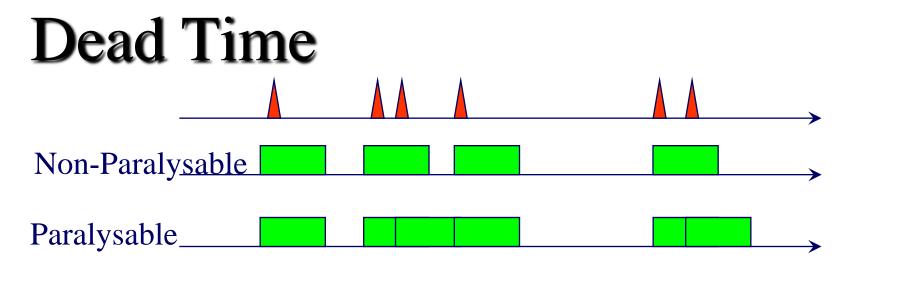
- Choosing optimum τ gives best noise performance but may not be fast enough
- We often have to sacrifice energy resolution for speed



- Shannon's Theorem and Nyquist Criterion
 - The highest frequency that can be 'measured' is twice the sampling frequency
- If the input is not band limited to frequencies less than $f_s/2$, then aliasing will occurs at frequencies $f \pm nf_s$
 - where f = signal frequency, fs = sampling frequency, n = integer
- If you have 100 μ m pixels, the ideal spatial resolution (PSF) > 200 μ m

Synchrotron Detectors

- A synchrotron source is used primarily when sensitivity is an issue
 - ♦ Signal too weak
 - Time resolution too poor
 - ♦ Sample too small
- More intensity can help this but...
- It places a major strain on detectors and Flux is a major issue for detectors!



 R_i =input rate, R_d =detected rate, τ dead time

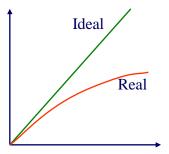
■ Non-paralysable

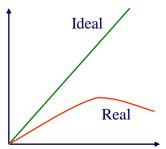
- Fraction of time detector is dead = $R_d \tau$
- Live time is therefore = $1 R_d \tau$
- Input rate = $R_i = R_d / (1 R_d \tau)$

Paralysable

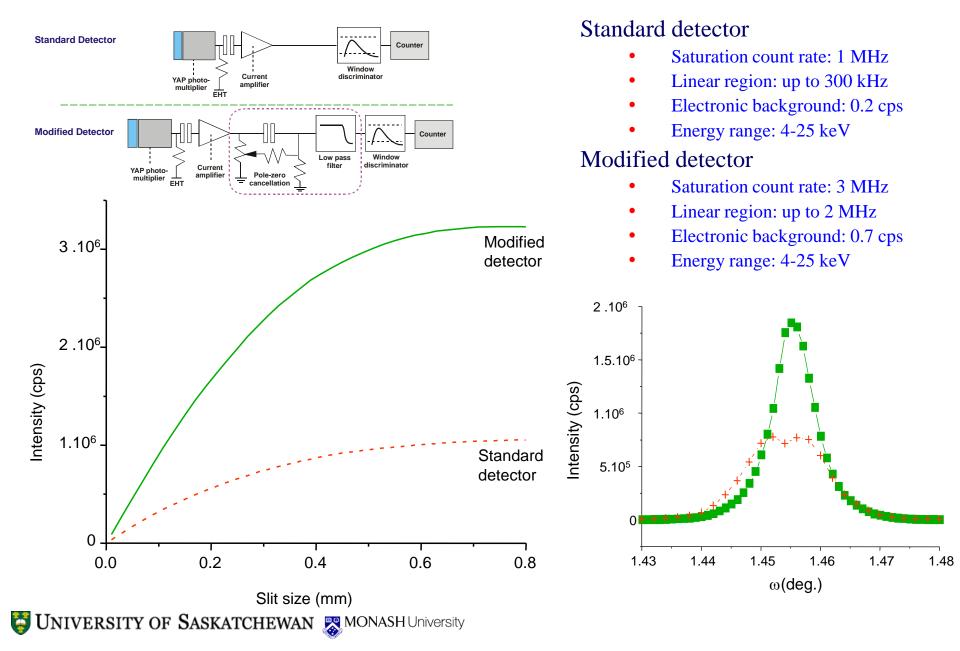
- R_d = Probability of getting no event within τ of an event $P(n,t) = \frac{e^{-R_i t} (R_i t)^n}{n}$
- Probability of n events in time t is

Detected rate $R_d = P(0, \tau) = R_i e^{-R_i \tau}$

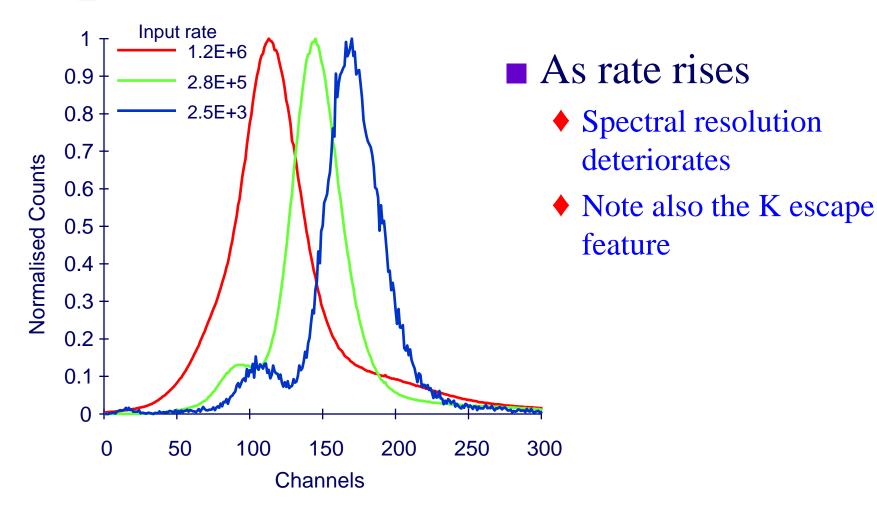




EDR Detector for Powder Diffraction



Spectral Peak Shift vs Rate



Detector Considerations

Intensity Measurement

- Uniformity across device
- Ageing, radiation damage
- Dynamic Range
- Linearity of Response
- ♦ Stability
- Spatial Measurement
 - Spatial Resolution
 - Spatial Distortion
 - ♦ Parallax

- Energy Measurement
 - Spectral Resolution
 - Linearity of Response
 - Uniformity of Response
 - ♦ Stability
- Time Measurement
 - Frame Rate
 - Photon Time Resolution
- Others
 - Size and weight
 - Cost

A Universal Specification?



Wombat

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Counting Statistics

- Photons are quantised and hence subject to probabilities
- The Poisson distribution expresses the probability of a number of events, k occurring relative to an expected number, n k = -n

$$P(n,k) = \frac{n^k e^{-n}}{k!}$$

- The mean of P(n, k) is n
- The variance of P(n, k) is n
- The standard deviation or error (noise) is \sqrt{n}
- If signal = n, then $SNR = n/\sqrt{n} = \sqrt{n}$
- As n increases, SNR improves

Performance Measure - DQE

Perfect detector

Real detector

$$SNR_{inc} = \sqrt{N_{inc}} \quad \therefore N_{inc} = SNR^{2}_{inc}$$
$$SNR_{Non-ideal} < \sqrt{N_{inc}}$$

Can define $N_{\mbox{\scriptsize photons}}$ that describes real SNR

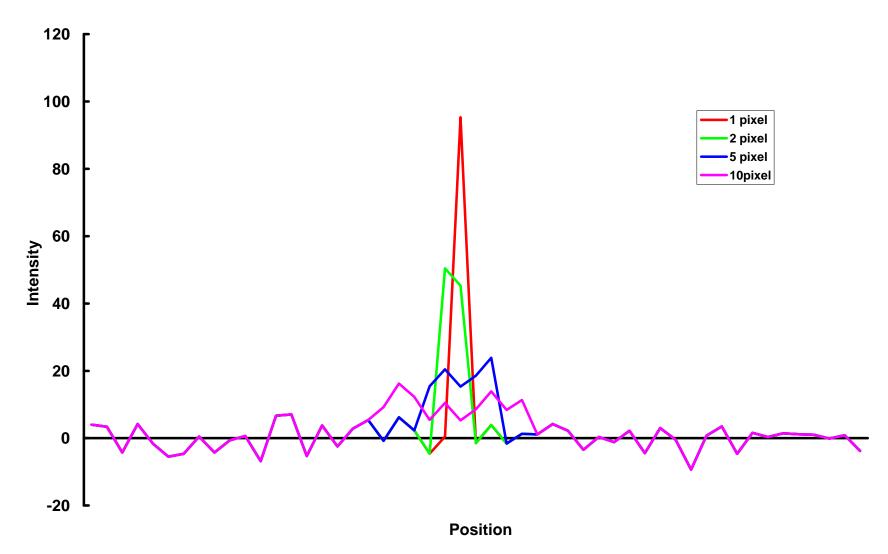
 $NEQ = SNR^2$ _{Non-ideal}

Ratio of this to N_{inc} is a measure of efficiency $DQE = \frac{NEQ}{N_{inc}} = \frac{SNR^2_{Non-ideal}}{SNR_{inc}^2}$

Note that DQE is f(spatial and spectral frequencies)

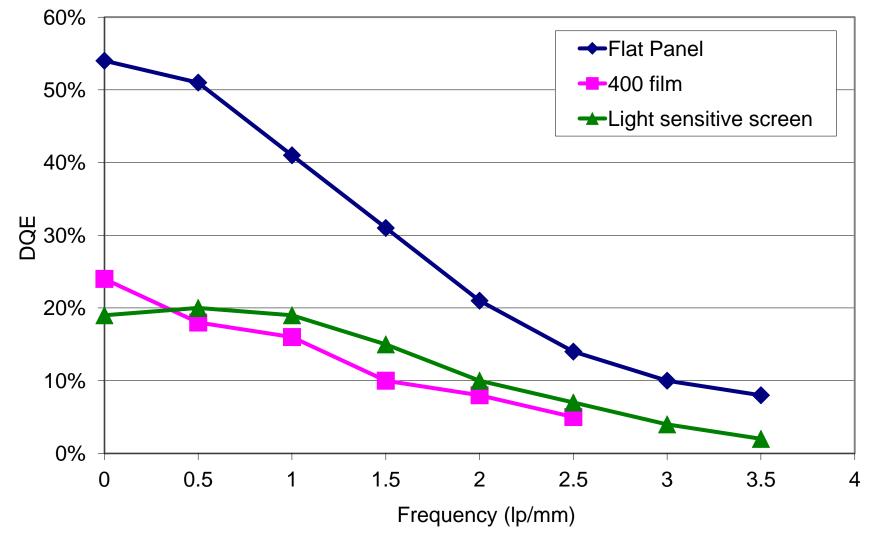
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Effect of Peak Width



DQE Comparison

DN-5 beam 2.6µGy



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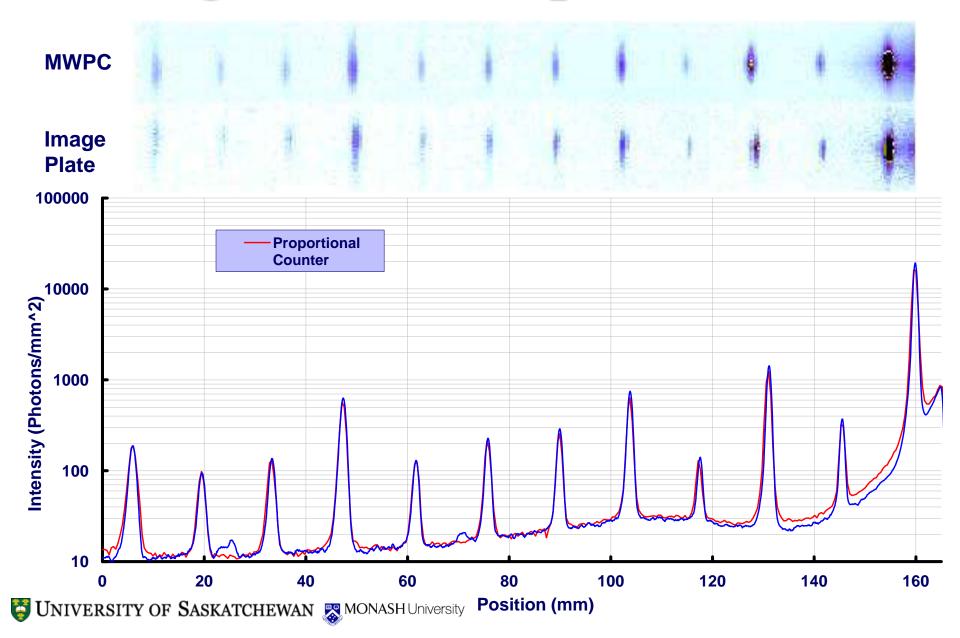
To Count or Not to Count



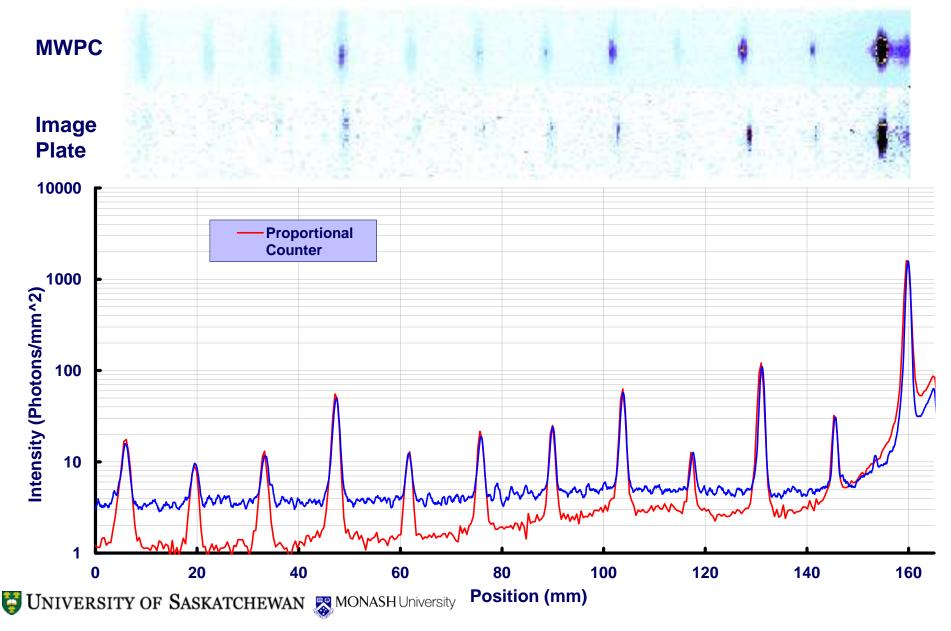
Tasmanian Devil



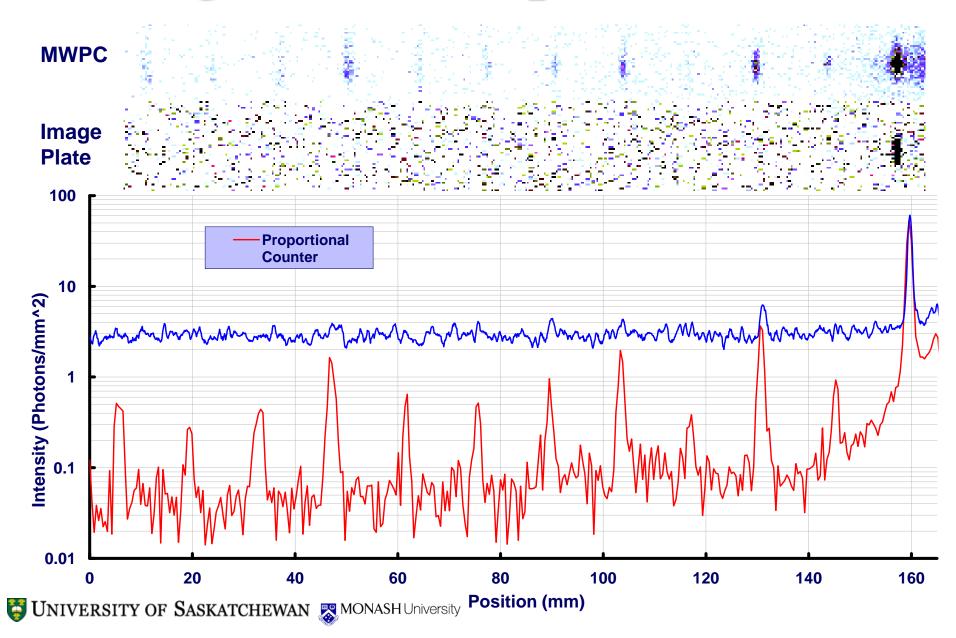
Collagen 100s Exposure



Collagen 10s Exposure



Collagen 0.3s Exposure



Cornell PAD (Integrating)

Rapid Framing Imager

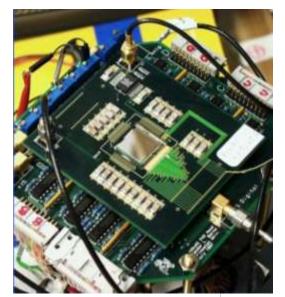
- ♦ 15×13.8mm² active area
- ♦ 150µm square pixel
- Storage for 8 frames

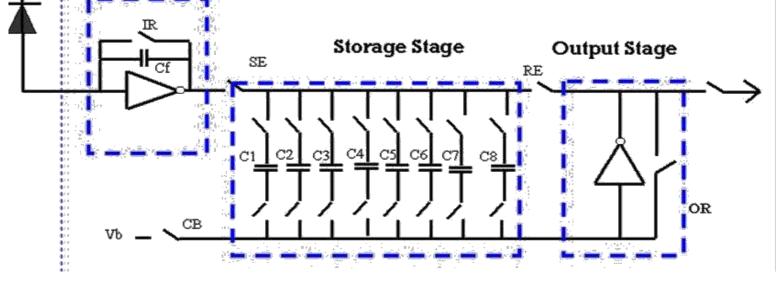
Input Stage

• Selectable T_{int} down to 1µs



+60V





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Sol Gruner, Cornell

Diesel Fuel Injection Movie

Injection

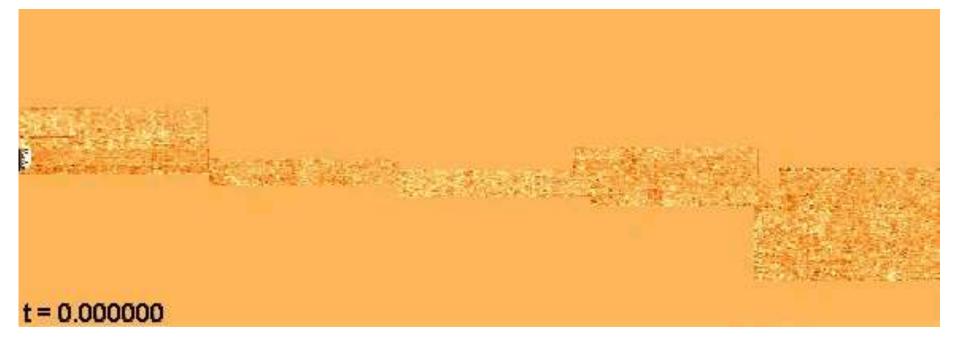
- Supersonic injection 1350psi Cerium added
- Chamber 1 atm SF_6
- 10⁸-10⁹ X-rays/s/pix (6keV)
- 1.1ms Pulse

Movie

- Length
- Frame length
- Dead time
- 168 frames (21 groups of 8)
- Average 20× to improve S/N
- Sequence

1.3ms

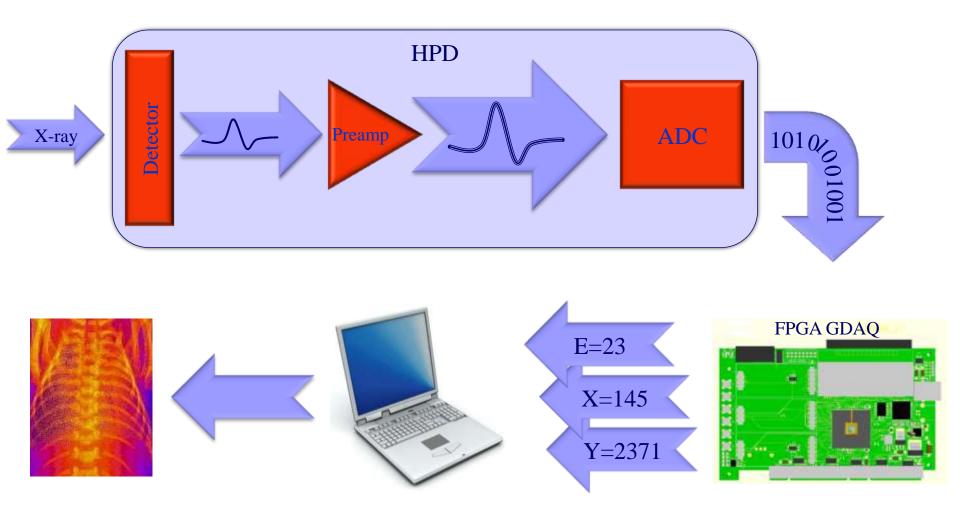
- 5.13µs
- $2.56\mu s$ / frame
- 5×10⁴ images



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A. MacPhee et al, Science (2002) 295, 1761-1763

Combine Imaging and Spectroscopy



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Pixel Array Detector

B

- A. Top electrode
- B. Pixellated semiconductor
- **C.** Collection electrodes

E

F

D. Bump bonds

C

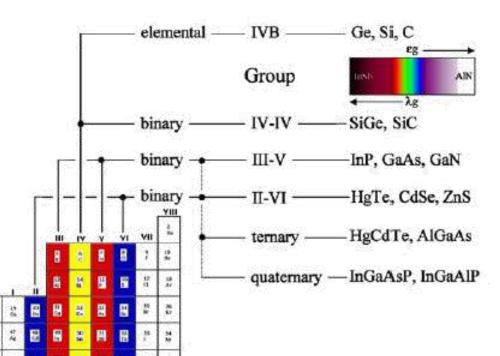
- **E.** Input electrode
- F. Pixellated ASIC

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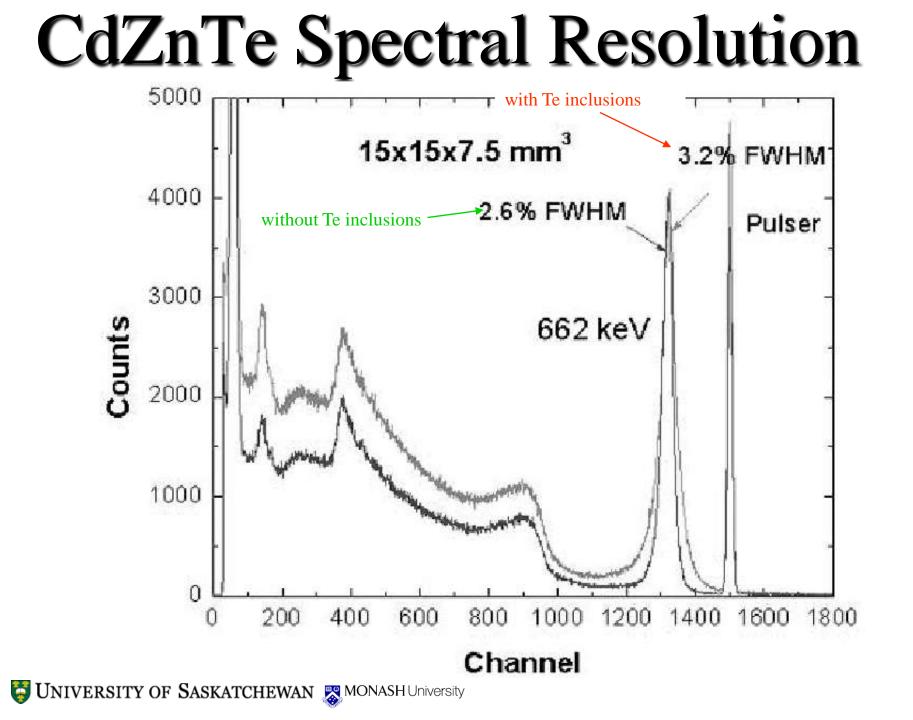
A

Available Compound Semiconductors

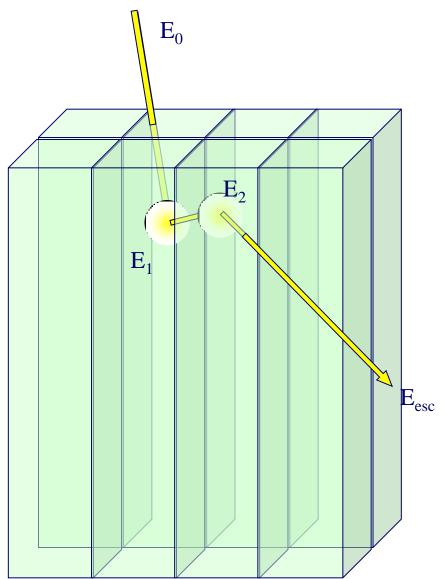
- Predominately CdZnTe, CdTe and GaAs.
- II-VI materials CdTe and CdZnTe cover a suitable range of band gaps:
 - 1.44 eV (CdTe), 1.57 eV (CdZnTe, 10% Zn), 1.64 eV (CdZnTe, 20% Zn)
- Resistivity of CdZnTe is higher than CdTe, hence lower dark current, higher spectroscopic resolution
- Poor hole transport requires electron-sensitive detectors



Paul Sellin, Surrey



The Problem of Multiple Scatters



• Need to measure E_0 • $E_0 = E_1 + E_2 + E_{esc}$

Must be able to detect multiple deposits as single event

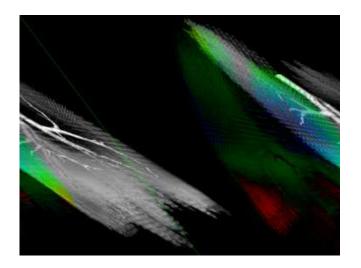
Must minimise E_{esc}



Other Issues

- In addition to detector performance metrics such as
 - ♦ Spatial resolution
 - Spectral resolution
 - ♦ Etc. etc.
- Consider other issues such as synchronisation.
- Many experiments require triggers or measurements of multiple parameters.

4D PIV



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S Dubsky, A Fouras et al

References

Delaney CFG and Finch EC

 Radiation detectors. Physical Principles and Applications, Clarendon Press, Oxford 1992, ISBN 0 19 853923 1

Knoll GE

- Radiation Detection and Measurement, John Wiley and Sons 2000
- Proceedings of the 7th International Conference on position sensitive detectors
 - Nuclear Instruments and Methods in Physics Research Volume 573, Issues 1-2, Pages 1-322
- IEEE Nuclear Science Symposia